

Searching for the Layered Structure of Space at the LHC

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Abstract

Alignment of the main energy fluxes along a straight line in a target plane has been observed in families of cosmic ray particles detected in the Pamir mountains. The fraction of events with alignment is statistically significant for families with superhigh energies and large numbers of hadrons. This can be interpreted as evidence for coplanar hard-scattering of secondary hadrons produced in the early stages of the atmospheric cascade development. This phenomenon can be described within the recently proposed “crystal world,” with latticized and anisotropic spatial dimensions. Planar events are expected to dominate particle collisions at a hard-scattering energy exceeding the scale Λ_3 at which space transitions from 3D \Rightarrow 2D. We study specific collider signatures that will test this hypothesis. We show that the energy-spectrum of Drell-Yan scattering is significantly modified in this framework. At the LHC, two jet and three jet events are necessarily planar, but four jet events can test the hypothesis. Accordingly, we study in a model-independent way the 5σ discovery reach of the ATLAS and CMS experiments for identifying four jets coplanarities. For the extreme scenario in which all $pp \rightarrow 4$ jet scattering processes become coplanar above Λ_3 , we show that with an integrated luminosity of $10(100) \text{ fb}^{-1}$ the LHC experiments have the potential to discover correlations between jets if $\Lambda_3 \lesssim 1.25(1.6) \text{ TeV}$.

I. INTRODUCTION

An intriguing alignment of gamma-hadron families (i.e., the outgoing high energy secondary particles from a single collision in the atmosphere) along a straight line in a target (transverse) plane has been observed with (lead and carbon) X -ray emulsion chambers (XREC's) in the Pamir mountains [1].¹ These families can be reconstructed by measuring the coordinates and the incident direction of each particle in the film emulsion. This allows determination of the total energy in gamma-rays and the total energy of hadrons release to gamma-rays. Recall that most of the hadrons in the family are pions and the average fraction of energy transferred by pions to the electromagnetic component is $\simeq 1/3$. All families in the experiment are classified by the value of the total energy observed in gamma-rays, $\sum E_\gamma$. The centers of the main energy fluxes deposited on the X -ray film (a.k.a. “subcores”) include halos of electromagnetic origin, gamma-ray clusters, single gamma-rays of high energy, and high energy hadrons. The criterion for alignment is given by the asymmetry parameter

$$\lambda_N = \frac{1}{N(N-1)(N-2)} \sum_{i \neq j \neq k} \cos 2\varphi_{ij}^k, \quad (1)$$

where N is the number of subcores and φ_{ij}^k is the angle between vectors issuing from the k -th subcore to the i -th and j -th subcores [2]. The parameter λ_N decreases from 1 (corresponding to N subcores disposed along a straight line) to $-1/(N-1)$ (corresponding to the isotropic case). Events are referred to as aligned if the N most energetic subcores satisfy $\lambda_N \geq \lambda_N^{\text{cut}}$. A common choice is $N = 4$ and $\lambda_N^{\text{cut}} = 0.8$.

The data have been collected at an altitude of 4400 m a.s.l., *i.e.*, at a depth of 594 g/cm² in the atmosphere. For low energy showers, 30 TeV $\lesssim \sum E_\gamma \lesssim 200$ TeV, the fraction of aligned events coincides with background expectation from fluctuations in cosmic ray cascade developments. However, for $\sum E_\gamma > 700$ TeV, the alignment phenomenon appears to be statistically significant [3]. Namely, the fraction (f) of aligned events is $f(\lambda_4 \geq 0.8) = 0.43 \pm 0.17$ (6 out of 14) in the Pb-XREC catalogue, and $f(\lambda_4 \geq 0.8) = 0.22 \pm 0.05$ (13 out of 59) in the C-XREC catalogue. The predominant part of the gamma-hadron families is produced by hadrons with energy $E_0 \gtrsim 10 \sum E_\gamma$, corresponding to interactions with a center-of-mass energy $\sqrt{s} \gtrsim 4$ TeV. Data analyses suggest that the production of most aligned groups occurs low above the chamber [2]. Thus, it is not completely surprising that the KASCADE Collaboration has found no evidence of this intricate phenomenon at sea level (~ 1000 g/cm²) [4].

Interestingly, the fraction of events with alignment registered in Fe-XREC's at Mt. Kanbala (in China) is also unexpectedly large [5]. For gamma-hadron families with energy $\sum E_\gamma \geq 500$ TeV the fraction of aligned events is $f(\lambda_3 \geq 0.8) = 0.5 \pm 0.3$ (3 out of 6). In

¹ The Pb-chambers are assembled of many sheets of lead (1 cm thick) interlaid with X -ray films. This provides a few interaction lengths for hadrons and a quasicolorimeter determination of the particle's energy. The C-chambers contain a 60 cm carbon layer covered on both sides by lead plates sandwiched with X -ray films. The carbon block provides a large cross section for hadron interaction, while the lead blocks are of minimal thickness allowing determination of particle energies. The total area of the chambers is few tens of square meters. Electron-photon cascades initiated by high energy hadrons and gamma-rays inside the XREC's produced dark spots whose sizes are proportional to the cascade energy deposited on the X -ray film.

addition, two events with $\sum E_\gamma \geq 1000$ TeV have been observed in stratospheric experiments [6]. Both events are highly aligned: (i) the so-called STRANA superfamily, detected by an emulsion chamber on board a Russian stratospheric balloon, has $\lambda_4 = 0.99$; (ii) the JF2af2 superfamily, detected by an emulsion chamber during a high-altitude flight of the supersonic aircraft Concord, has $\lambda_4 = 0.998$. It is worth noting that stratospheric experiments record the alignment of particles, whereas mountain-based facilities register the alignment of the main fluxes of energy originated by these particles on a target plane.

The strong collinearity of shower cores has been interpreted as a tendency for coplanar scattering and quasiscaling spectrum of secondary particles in the fragmentation region [7]. If the aligned phenomenon observed in cosmic ray showers is not a statistical fluctuation, then events with unusual topology may be produced at the Large Hadron Collider (LHC). In this paper we carry out a systematic study of $pp \rightarrow 4$ jet scattering processes to establish the sensitivity of LHC experiments to such planar-shape topology.

The analysis technique described herein constitute an entirely general approach to search for planar scattering at the LHC. Firstly, we generate Standard Model (SM) QCD events that contain 4 light jets and 2 b -quark jets and 2 light jets utilizing the ALPGEN Monte Carlo code [8]. Next, the event shape variable is classified according to the standard aplanarity parameter, A_p [9]. Namely, we define the signal region as $A_p < A_p^{\text{cut}}$ and the control region with $A_p > A_p^{\text{cut}}$, varying the cut among these three choices: $A_p^{\text{cut}} = 0.1, 0.05, 0.01$. After that, we extract the ratio of bi-planar to planar events (N_b/N_p) and calculate the N_b/N_p uncertainty based on Poisson statistics for a given luminosity L . Finally, we calculate the required luminosity to obtain 5σ , 3σ , and 95% C.L. away from various values of N_b/N_p . Though the search technique is agnostic regarding the hypothetical physics underpinning the planar configuration of events, an observation could have bearing on the recently proposed idea that spatial dimensions collapse at short distance, shutting off one-by-one with rising energy [10].

The layout of the paper is as follows. In Sec. II we study the LHC sensitivity to coplanar particle escape. We quantify signal and background rates of 4 jet events and show that future measurements of the aplanarity distribution of multi-jet events can provide a potent method for exposing a dimensional reduction of phase space. In Sec. III we associate the dimensional reduction in momentum space to a reduction in spatial dimensions. We assume that space at its fundamental level is an anisotropic lattice [10]. This idea that spatial dimensions effectively reduce with increasing energy directly contrasts with field/string theories in continuous spacetime dimensions, where dimensionality increases with a rise in energy. In Sec. IV we present a phenomenological analysis of the Drell-Yan scattering processes [11] for a $3D \Rightarrow 2D$ crossover and we show that Tevatron data is insufficient to constrain the model for a dimensional reduction above 1 TeV. We summarize our conclusions in Sec. V.

II. SENSITIVITY OF THE LHC EXPERIMENTS TO 4 PLANAR-JET EVENTS

In this section we estimate the LHC sensitivity to coplanar events and corresponding discovery reach at $\sqrt{s} = 14$ TeV using $2 \rightarrow 4$ scattering processes. Such scattering processes involve multiple virtual particles. We define Λ_3 as the energy-scale of the onset of new physics. We assume that when the momentum transfer Q ($Q^2 = -\hat{t}$) in each of the propagators is comparable with Λ_3 , a growing fraction of the jets are produced in one plane in their center-of-mass frame. This coplanarity is drastically different from the usual topology of 4D scattering, where the four outgoing partons are in general acoplanar. For simplicity,

we consider three values of the coplanarity fraction at Λ_3 . We investigate the fractions 30%, 50%, and 100%, utilizing simulated QCD events (with ALPGEN [8]) containing 4 jets (either 4 light jets or 2 b -quark jets and 2 light jets). While a fraction of the simulated events is planar, the entire sample is generally bi-planar. Given the limited number of events containing four high transverse momentum (p_T) jets, the probe of truly planar events is statistics limited. Therefore, we look to compare the SM QCD prediction of bi-planar jets with a purely planar sample by determining how many events are required for the observation of *aplanar* events to be significant.

For a given propagator scale, Λ_3 , we accept events that have 4 jets that pass the following acceptance cuts:

$$p_{T,j} \geq \frac{1}{2}\Lambda_3 \text{ GeV}, \quad \eta_j < 2.5, \quad \Delta R_{j,j} > 0.4, \quad (2)$$

where η_j is the j jet pseudorapidity, and $\Delta R_{j,j}$ is the separation in the azimuthal angle (ϕ) - pseudorapidity (η) plane among jets:

$$\Delta R_{j1,j2} = \sqrt{(\phi_{j1} - \phi_{j2})^2 + (\eta_{j1} - \eta_{j2})^2}. \quad (3)$$

In addition, to reject events that do not have a hard momentum scale, $Q \sim \Lambda_3$, we require that the invariant mass of any two pairs of jets satisfy

$$M_{jj} > \Lambda_3, \quad (4)$$

and that for any two pairs of jets, the transverse momentum of a jet relative to the boost axis of the jet pair is

$$k_T > \frac{\Lambda_3}{2}. \quad (5)$$

Finally, we model detector resolution effects by smearing the final state jet energy according to:

$$\frac{\Delta E}{E} = \frac{0.50}{\sqrt{E/\text{GeV}}} \oplus 0.03. \quad (6)$$

The event shape variable, aplanarity, can be calculated based on the constructed momentum tensor

$$M_{ab} = \frac{\sum_i k_{ia} k_{ib}}{\sum_i k_i^2} \quad (7)$$

where i runs over the 4 jets in each event, with all momenta measured in the c.m. The aplanarity is defined as [9]

$$A_p = \frac{3}{2}Q_1, \quad (8)$$

where Q_1 is the smallest normalized eigenvalue of the momentum tensor, giving A_p a maximum value of $\frac{1}{2}$. Therefore, planar or collinear events possess $A_p \sim 0$ values, while more 3D events approach the maximum value, $A_p = \frac{1}{2}$.

Now, we define events as exhibiting a planar topology by having the aplanarity below some cut value, $A_p < A_p^{\text{cut}}$. In Fig. 1.a, we show A_p for a variety of jet p_T cuts for the LHC with 100 fb^{-1} of integrated luminosity. To facilitate distinguishing planar events from the usual SM QCD bi-planar events, we define the variable

$$\rho_O = N_b/N_p, \quad (9)$$

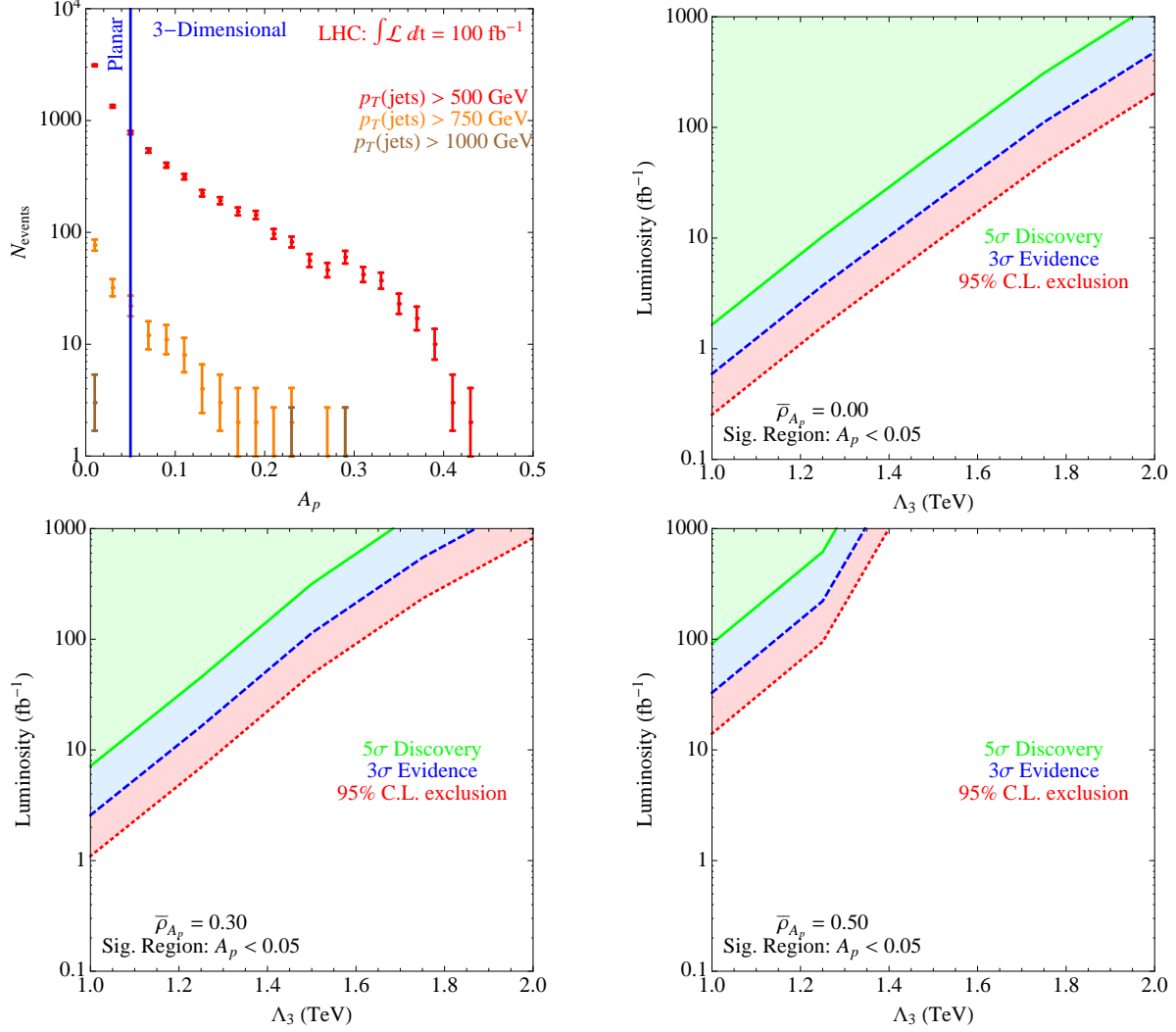


FIG. 1: (a) Aplanarity of SM QCD 4-jet events for 100 fb $^{-1}$ and jet p_T acceptance cuts (from above) of 500, 750 and 1000 GeV. We take events with $A_p < 0.05$ as being planar. (b) Reach of Λ_3 at the LHC based on the aplanarity event shape variable. The 5 σ discovery is indicated by a solid line, the 3 σ evidence by a dashed line, and the 95% CL exclusion by a dotted line.

the ratio of 3D bi-planar events to planar events using the observable \mathcal{O} as a discriminator for planar and bi-planar events. For pure planar events, this ratio should vanish, while for 3D SM events, it is generally nonzero. For a given p_T cut, we can identify at what luminosity it is statistically different from a ratio the model predicts, i.e. $\bar{\rho}_{\mathcal{O}}$. We vary the value of A_p^{cut} and find that $A_p^{\text{cut}} = 0.05$ gives the maximum sensitivity for $\bar{\rho}_{A_p} = 0$, denoted by the vertical line in Fig. 1.a. In Fig. 1.(b-d), we show the reach the LHC may have for a given luminosity for planar events, recast into the propagator scale, Λ_3 for values of $\bar{\rho}_{A_p}$. With 10(100) fb $^{-1}$, the LHC may discover propagators of scale $\Lambda_3 \approx 1.25(1.6)$ TeV with $\bar{\rho}_{A_p} = 0$. The reach is degraded as $\bar{\rho}_{A_p}$ increases such that at $\bar{\rho}_{A_p} = 0.5$, 5 σ discovery at 100 fb $^{-1}$ may be possible only for $\Lambda_3 = 1$ TeV. This strong dependence on $\bar{\rho}_{A_p}$ is due to the 3D region in Fig. 1a having fewer events than the planar region.

In addition to the aplanarity, one can construct other variables to test for planar events.

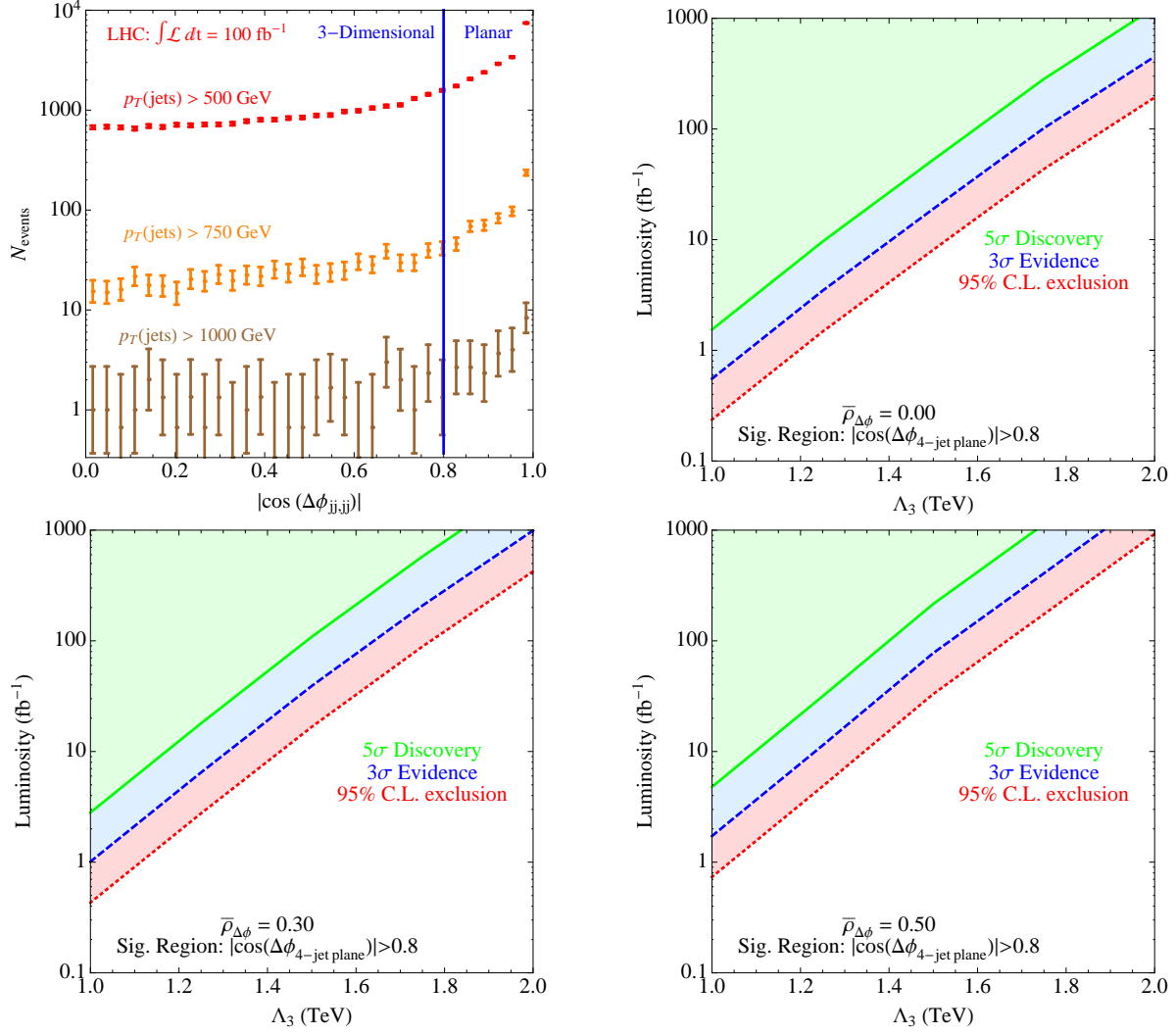


FIG. 2: (a) The $|\cos(\Delta\phi_{jj,jj})|$ for 100 fb^{-1} and jet p_T acceptance cuts (from above) of 500, 750 and 1000 GeV, where $\Delta\phi_{jj,jj} \equiv \Delta\phi_{4\text{-jet plane}}$ is the angle between the two planes defined by the respective jets combinations. There are a total of 3 combinations, and each are included. We define the signal region as $|\cos(\Delta\phi_{4\text{-jet plane}})| > 0.8$. (b) Reach of Λ_3 at the LHC based on $|\cos(\Delta\phi_{4\text{-jet plane}})|$. The 5 σ discovery is indicated by a solid line, the 3 σ evidence by a dashed line, and the 95% CL exclusion by a dotted line.

For instance, two planes may be defined as going through each jet pair. The subsequent azimuthal angle, $\Delta\phi_{jj,jj}$, between these two planes should be zero for all combinations of jet pairings for planar events, and nonzero otherwise. When using this alternate variable and placing a cut as done for the A_p analysis above (see Fig. 2.a), we arrive at a very similar reach for $\bar{\rho}_{\Delta\phi} = 0$, shown in Fig. 2.b; other $\bar{\rho}_{\Delta\phi} \neq 0$ cases are shown in Fig. 3(c,d). The weaker dependence of the reach on the nonzero value of $\bar{\rho}_{\Delta\phi}$ than in the A_p case is due to the smaller ρ_{A_p} value in the SM. Thus, making it easier to statistically distinguish between the bi-planar SM and planar events.²

² Note, however, that the values $\bar{\rho}_{A_p}$ and $\bar{\rho}_{\Delta\phi}$ are not necessarily equivalent.

An approach for production of highly coplanar multi-jet events, which exploits semi-hard QCD processes with a high transverse momentum transfer, has been discussed in [12]. This model is successful in explaining coplanar 3-jet events in cosmic ray data. The leading particle provides one jet and the collinear singularity of QCD correlates the second and third jets, therefore producing a roughly coplanar 3-jet event in the lab system. The enhanced amplitude due to collinear gluon emission may explain, for example, the two jets “ridge” phenomenon recently observed at the LHC in pp collisions [13] and previously observed at RHIC in heavy ion collisions [14]. However, QCD collinearity cannot explain $N \geq 4$ jet events in cosmic ray physics, which corresponds to $N \geq 3$ jet events in the center-of-mass system (which is also the lab system) at the LHC. In the remainder of this paper we will discuss a model where multi-jet events are coplanar for all values of N .

III. THE CRYSTAL WORLD

Motivated by condensed matter systems, some of us recently proposed spacetime may be an ordered lattice structure that becomes anisotropic at very small distances [10]. The proposed set up, shown in Fig. 3, resembles that of dimensional crossover in layered strongly correlated metals [15]. These materials have an insulating character in the direction perpendicular to the layers at high temperatures but become metal-like at low temperatures, whereas transport parallel to the layers remains metallic over the whole temperature range. The analogy which we adopt is to replace the temperature variable in the materials system with short-distance “virtuality” in the parton scattering processes. We further assume that the lattice orientation is randomized on a scale sufficiently small to avoid any preferred direction in space on the macroscopic scale. On the small scale, there will be a preferred direction given by the local lattice orientation. Therefore, hard scattering processes can resolve the lattice spacings that separate “conducting space” from the “crystal world” of insulating space. On the other hand “macroscopic” objects like beam protons effectively see a spacetime continuum.

It is of interest to explore the consequences of this extreme viewpoint (and let experiment be the arbiter). We are not aware of any data that would rule out this conjecture. We emphasize that this conjecture is radically different from the increase of dimensions at small spatial scales that are postulated in string theories and in many modern field theories. In contrast with our conjecture the effective number of dimensions decreases as partons probe smaller scales. With the reduction in spatial dimension, phase space is reduced, the cross-sections are reduced, and multi-jet final states are necessarily coplanar (for $3 \rightarrow 2$). With an increase in spatial dimension, phase space is increased, cross-sections increase, and multi-jet final states fill the three spatial dimensions but also lose energy and multiplicity into the extra dimensions.

As demonstrated in Sec. II for the coplanarity, striking collider signals would be observed at the LHC if $\Lambda_3 \sim 1/L_3 \sim 1$ TeV. Many related aspects of beyond the SM field theory will occur at high-energy: the transition of the renormalizable SM to a super-renormalizable field theory, modification in the evolution properties of parton distribution functions, running coupling constants, running anomalous dimensions of operators, and so on. A study of these new effects requires a more specific model. Such a study is beyond the scope of this paper. If the LHC provides indications for the correctness of our conjecture, then it becomes necessary and even imperative to explore the new physics in detail.

In the next section we address an obvious probe of new physics, namely Drell-Yan pro-

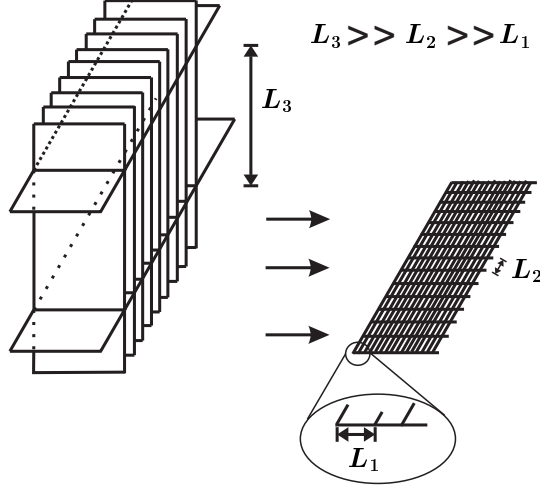


FIG. 3: Ordered lattice. The fundamental quantization scale of space is indicated by L_1 . Space structure is $1D$ on scales much shorter than L_2 , while it appears effectively $2D$ on scales much larger than L_2 but much shorter than L_3 . At scales much larger than L_3 , the structure appears effectively $3D$ [10].

duction of lepton pairs. Above $\sqrt{s} \sim \Lambda_3$, the reduction of phase space is expected to reduce the Drell-yan cross-section, We show that present Tevatron data does not impact on the conjecture, but that future LHC data will provide a new reach in Λ_3 .

IV. DRELL-YAN MEETS THE CRYSTAL WORLD

In this section we assess the effect on Drell-Yan cross sections at Tevatron and at LHC of the restriction to coplanarity of scattering above the lattice energy $\sim \Lambda_3 \sim 1$ TeV. In order to effectuate this, we will need to hypothesize an algorithm for a smooth continuation between $(2+1)D$ to $(3+1)D$.

In a standard manner, we define the cross section as the probability transition rate divided by the flux, and obtain

$$\begin{aligned} d\sigma_D &= \frac{1}{(2\pi)^{D-2}} \frac{1}{16\hat{s}} \frac{p_f^*}{p_i^*} |\mathcal{M}_D|^2 d\Omega_{D-1}^* |\vec{p}_f^*|^{D-4} \\ &= \left(\frac{|\mathcal{M}_D|}{8\pi\sqrt{\hat{s}}} \right)^2 \left(\frac{p_f^*}{p_i^*} \right) d\Omega_{D-1}^* \left(\frac{|\vec{p}_f^*|}{2\pi} \right)^{D-4}, \end{aligned} \quad (10)$$

where \mathcal{M}_D is the Feynman amplitude invariant in D dimensions; p_i^* and p_f^* are the initial and final center-of-mass momenta $\sim \sqrt{\hat{s}}/2$; and $d\Omega_{D-1}^*$ is the solid angle aperture in the center of mass. Thus

$$\sigma_4 \propto \mathcal{M}_4^2 / \hat{s}, \quad (11)$$

whereas for $D = 3$

$$\sigma_3 \propto \mathcal{M}_3^2 / \hat{s}^{3/2}. \quad (12)$$

(Note that since $\mathcal{M}_3 \sim E^1$, the “cross section” σ_3 in two spatial dimensions is a straight line with dimension length, as expected.)

To assess the impact on Drell-Yan, it is of critical importance to modify σ_3 in order that it gain entry as a σ_4 in the parton model calculation of $pp(\bar{p}) \rightarrow \ell^+\ell^- + X$. This entails a change in dimension for the cross sections, and thus it *cannot* occur in the probability transition rate, which has the same dimension (E^1) for any D . Thus it must reside in the incoming flux [16]. Assuming that this modification is purely of geometric origin, without energy dependence, we insert a factor of Λ_3 to adjust dimensions and obtain the following prescription for the transition to $D = 3$:

$$\begin{aligned}\sigma_4^{\text{eff}} &\sim \sigma_3/\Lambda_3 \\ &\sim (\mathcal{M}_3^2/\Lambda_3) / \hat{s}^{3/2}\end{aligned}\tag{13}$$

At this point we must deal with the matrix element \mathcal{M}_3 . There are essentially two choices: (i) one can assume total ignorance of the dynamics, (seeing as it may involve the interaction of particles with a space time lattice), and simply work on dimensional grounds. This approach is decidedly unsatisfactory, since it would permit arbitrary powers of $\sqrt{\hat{s}}/\Lambda_3$, vitiating any hope of making contact with quantitative experimental findings; or (ii) we can assume that \mathcal{M}_3 can be calculated from a 3 dimensional version of QED. In case (ii), the operational Lagrangian is

$$\mathcal{L}_{\text{QED}_3} = i\bar{\psi}\gamma^\mu(\partial_\mu - ie_3 A_\mu)\psi - \frac{1}{4}F^{\mu\nu}F_{\mu\nu},\tag{14}$$

where $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ is the field strength, e_3 is the coupling, and ψ is a four-component spinor with the corresponding four-dimensional representation of the Dirac algebra

$$\gamma^0 = \begin{pmatrix} \sigma^3 & 0 \\ 0 & -\sigma^3 \end{pmatrix}, \quad \gamma^1 = \begin{pmatrix} i\sigma^1 & 0 \\ 0 & -i\sigma^1 \end{pmatrix}, \quad \gamma^2 = \begin{pmatrix} i\sigma^2 & 0 \\ 0 & -i\sigma^2 \end{pmatrix},\tag{15}$$

with σ^i the Pauli matrices [17].

Then the calculation of Drell-Yan proceeds the same manner as in $D = 4$, with the result

$$\mathcal{M}_3 \sim e_3^2 \hat{s}/Q^2 \sim e_3^2.\tag{16}$$

where the factor \hat{s} comes from the Dirac spinor normalization and the Q^2 from the virtual photon (or Z) propagator. An engineering dimensional analysis of (14) shows $e_3^2 \sim E^1$, so that indeed $\mathcal{M}_3 \sim e_3^2 \sim \text{const} \sim \Lambda_3$.³ In this case, Eq.(13) reads

$$\sigma_4^{\text{eff}} \sim \Lambda_3/\hat{s}^{3/2}\tag{17}$$

for $\sqrt{\hat{s}} > \Lambda_3$. More explicitly, the working algorithm for examining effects of the transition to coplanarity would be the following:

$$\begin{aligned}\sigma^{\text{DY}} &= \sigma_{\text{SM}}^{\text{DY}} && \sqrt{\hat{s}} \ll \Lambda_3 \\ &= \sigma_{\text{SM}}^{\text{DY}} (\Lambda_3/\sqrt{\hat{s}}) && \sqrt{\hat{s}} \gg \Lambda_3,\end{aligned}\tag{18}$$

where $\sigma_{\text{SM}}^{\text{DY}}$ is the Drell-Yan parton-parton cross section as calculated in the SM. For convenience we use as an interpolating function

$$\sigma^{\text{DY}} = \sigma_{\text{SM}}^{\text{DY}} \left(1 - e^{-\Lambda_3/\sqrt{\hat{s}}}\right),\tag{19}$$

³ In a more general situation e_3 would be energy (and dimension) dependent, receiving contributions from radiative corrections [18].

where $e^{-\Lambda_3/\sqrt{s}}$ parameterizes the the probability to remain confined to the two dimensional space. The quantity Λ_3 may also be regarded as characteristic of a potential barrier for phase transitions from 3 to 2 spatial dimensions, with $e^{-\Lambda_3/\sqrt{s}}$ representing the probability of tunneling between these states. This is perhaps reminiscent of the the factor $\sim \mathcal{E}^2 e^{-\pi m^2/e|\mathcal{E}|}$ the probability per unit time per unit volume for creating an e^+e^- pair in a constant electric field \mathcal{E} [19]. To obtain the total cross section for the process $pp(\bar{p}) \rightarrow \ell^+\ell^- + X$ we have to convolute (18) with the parton distribution functions.

We do not at present have a controlled calculation coming from a well-defined formalism for the parton distribution with particle momenta oblique to the lattice layer. A crude approximation that conserves 4-momentum among the particles is to assume the beam axis is aligned with the lattice layer. This hadron alignment in turn means that the partons in the infinite momentum frame are aligned with the lattice. In reality, it is only a projection of the parton momentum that is aligned with the lattice. Each of the two partons oblique to the lattice will presumably undergo an inelastic scattering with the lattice plane in advance of their mutual hard scattering.⁴ This new mechanism of parton energy loss is somewhat analogous to an initial state radiation off the partons beyond the physical radiation included in factorization theorems.

A question that immediately arises is whether there are constraints on this model from Tevatron data. The D0 Collaboration reported the most recent study of the dielectron invariant mass spectrum analyzing 5.4 fb^{-1} of data collected at $\sqrt{s} = 1.96 \text{ TeV}$ [20]. There are four events observed in the energy bin $600 \text{ GeV} - 800 \text{ GeV}$. Since there is negligible background from other SM processes, the 68.27%CL spread in the Poisson signal mean is (2.34, 6.78) events [21]. At this point it is worth recalling that the Drell-Yan $d\sigma^{\text{DY}}/dQ$ measures lepton pair spectra at parton collision energy Q . Therefore, from Eq. (19) it is straightforward to verify that a dimensional reduction at $\Lambda_3 = 1 \text{ TeV}$ (which predicts $\simeq 3$ events in the $600 \text{ GeV} - 800 \text{ GeV}$ energy bin) is consistent with Tevatron data at the 1σ level. A similar analysis follows from the invariant mass spectrum of dimuon data reported by the CDF Collaboration [22]. Interestingly, the $\Lambda_3 \sim 1 \text{ TeV}$ region will be tested by the early LHC run at $\sqrt{s} = 7 \text{ TeV}$.

V. CONCLUSIONS

In the first part of this paper we have presented a complete model independent study to search for planar events in 4 jet final states. The only free parameter of our analysis is the characteristic energy scale for the onset of coplanarity, Λ_3 . For the extreme scenario in which all $pp \rightarrow 4$ jet scattering processes become coplanar above Λ_3 , we have shown that with an integrated luminosity of $10(100) \text{ fb}^{-1}$ the LHC experiments have the potential to discover correlations between jets if $\Lambda_3 \lesssim 1.25(1.6) \text{ TeV}$.

In the second part of this paper we have discussed some aspects of the recently proposed [10] latticized spatial dimensions, with different characteristic lattice spacings in each dimension. We have studied specific collider signatures that will probe this idea. In particular, we have shown that the predicted energy spectrum of Drell-Yan scattering is significantly modified in this model. Remarkably, the anisotropic crystal world yields planar events when

⁴ The single parton is assumed to coherently, elastically scatter from the lattice, thereby altering its refractive index (velocity) in a small way.

the energies of hard scatterings exceed the scale at which space transitions from 3D \rightleftharpoons 2D. Therefore, four jet events at the LHC will exhibit striking planar alignment (if the parton-momentum transfer Q exceeds the energy scale Λ_3 of the lattice). Jets with this strong azimuthal anisotropy may have been already observed by the Pamir Collaboration: the effect known as alignment, which cannot be explained by conventional physics.

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